Computational & Multiscale Mechanics of Materials

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April 2019 CM3 research projects

Direct links

- Mean-Field-Homogenization for Elasto-Visco-Plastic Composites
- <u>Micro-structural simulation of fiber-reinforced highly crosslinked epoxy</u>
- Stochastic Homogenization of Composite Materials
- Bayesian identification of stochastic MFH model parameters
- Non-Local Damage Mean-Field-Homogenization
- Boundary conditions and tangent operator in multi-physics FE²
- <u>Computational Homogenization For Cellular Materials</u>
- <u>Stochastic 3-Scale Models for Polycrystalline Materials</u>
- DG-Based Fracture
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 - DG-Based Dynamic Fracture
 - DG-Based Damage elastic damage to crack transition
- Non-local Gurson damage model to crack transition
- <u>Stochastic Multi-Scale Fracture of Polycrystalline Films</u>
- <u>Smart Composite Materials Shape Memory Effects</u>
- Multi-Scale Modeling of Nano-Crystal Grain Boundary Sliding
- <u>Stochastic Multi-Scale Model to Predict MEMS Stiction</u>



Beginning

Computational & Multiscale Mechanics of Materials





Mean-Field-Homogenization for Elasto-Visco-Plastic Composites

SIMUCOMP The research has been funded by the Walloon Region under the agreement no 1017232 (CT-EUC 2010-10-12) in the context of the ERA-NET +, Matera + framework.

The authors gratefully acknowledge the financial support from F.R.S-F.N.R.S. under the project number PDR T.1015.14 STOMMMAC The research has been funded by the Walloon Region under the agreement no 1410246-STOMMMAC (CT-INT 2013-03-28) in the context of M-ERA.NET Joint Call 2014.



April 2019 CM3 research projects

• Multi-scale modeling

- 2 problems are solved concurrently
 - The macro-scale problem
 - The meso-scale problem (on a meso-scale Volume Element)





$$L_{\text{macro}} >> L_{\text{VE}} >> L_{\text{micro}}$$

For accuracy: Size of the mesoscale volume element smaller than the characteristic length of the macro-scale loading To be statistically representative: Size of the meso-scale volume element larger than the characteristic length of the microstructure





- Remove residual stress in matrix
 - Or use second moment estimates

 $\Delta \boldsymbol{\epsilon}_{I}^{r} = \boldsymbol{B}^{\varepsilon} (I, \bar{\boldsymbol{C}}_{0}^{S0}, \bar{\boldsymbol{C}}_{I}^{Sr}) : \Delta \boldsymbol{\epsilon}_{0}^{r} \qquad \& \boldsymbol{\sigma}_{0} = \bar{\boldsymbol{C}}_{0}^{S0} : \Delta \boldsymbol{\epsilon}_{0}^{r}$

Beginning

- Incremental-secant mean-fieldhomogenization
 - Stress tensor (2 forms)

 $\begin{cases} \boldsymbol{\sigma}_{I/0} = \boldsymbol{\sigma}_{I/0}^{res} + \bar{\boldsymbol{C}}_{I/0}^{Sr} : \Delta \boldsymbol{\varepsilon}_{I/0}^{r} \\ \boldsymbol{\sigma}_{I/0} = \bar{\boldsymbol{C}}_{I/0}^{S0} : \Delta \boldsymbol{\varepsilon}_{I/0}^{r} \end{cases}$

- Radial return direction toward residual stress
 - First order approximation in the strain increment (and not in the total strain)
 - Exact for the zero-incremental-secant method
- The secant operators are naturally isotropic

$$\begin{cases} \bar{\mathbf{C}}^{\mathrm{Sr}} = 3\kappa^{\mathrm{el}}\mathbf{I}^{\mathrm{vol}} + 2\left(\mu^{\mathrm{el}} - 3\frac{{\mu^{\mathrm{el}}}^2\Delta p}{\left(\boldsymbol{\sigma}_{n+1} - \boldsymbol{\sigma}_n^{\mathrm{res}}\right)^{\mathrm{eq}}}\right)\mathbf{I}^{\mathrm{vol}} \\ \bar{\mathbf{C}}^{\mathrm{S0}} = 3\kappa^{\mathrm{el}}\mathbf{I}^{\mathrm{vol}} + 2\left(\mu^{\mathrm{el}} - 3\frac{{\mu^{\mathrm{el}}}^2\Delta p}{\boldsymbol{\sigma}_{n+1}^{\mathrm{eq}}}\right)\mathbf{I}^{\mathrm{vol}} \end{cases}$$





Beginning

- Incremental-secant mean-field-homogenization
 - Second-statistical moment estimation of the von Mises stress
 - First statistical moment (mean value) not fully representative

$$\overline{\sigma}_{I/0}^{eq} = \sqrt{\frac{3}{2}} \overline{\sigma}_{I/0}^{dev} : \overline{\sigma}_{I/0}^{dev}$$

• Use second statistical moment estimations to define the yield surface



- Non-proportional loading
 - Spherical inclusions
 - 17 % volume fraction
 - Elastic
 - Elastic-perfectly-plastic matrix





- Elasto-visco-plasticity
 - Elasto-visco-plastic short fibres
 - Spherical
 - 30 % volume fraction
 - Elasto-visco-plastic matrix





• Extension to finite deformations

- Formulate everything in terms of elastic left Cauchy-Green tensor





Beginning





• SIMUCOMP ERA-NET project (incremental secant MFH)

- e-Xstream, CENAERO, ULiège (Belgium)
- IMDEA Materials (Spain)
- CRP Henri-Tudor (Luxemburg)
- PDR T.1015.14 project (MFH with second-order moments)
 - ULiège, UCL (Belgium)

• STOMMMAC M.ERA-NET project (MFH for elasto-visco-plastic composites)

- e-Xstream, ULiège (Belgium)
- BATZ (Spain)
- JKU, AC (Austria)
- U Luxembourg (Luxemburg)
- Publications (doi)
 - <u>10.1016/j.mechmat.2017.08.006</u>
 - <u>10.1080/14786435.2015.1087653</u>
 - <u>10.1016/j.ijplas.2013.06.006</u>
 - <u>10.1016/j.cma.2018.12.007</u>



Beainnina

Computational & Multiscale Mechanics of Materials





Micro-structural characterization and simulation of fiber-reinforced highly crosslinked epoxy

The authors gratefully acknowledge the nancial support from F.R.S-F.N.R.S. under the project number PDR T.1015.14



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Resin behavior (experiments UCL)

- Viscoelasto-Viscoplaticity
- Saturated softening

2.5

2

1.5

 0.5^{1}

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-6

 $\sigma^{\rm eq}/\sigma_c$

- Asymmetry tension-compression
- Pressure-dependent yield

To used in micro-structural analysis

- Behavior in composite is different
- Introduce a length-scale effect

 $\alpha = 2$ (Paraboloid)

Exp. Lesser 1997

Exp. Hinde 2005

Exp. Sauer 1977

-4

-2

 p/σ_c

α=3.5

α=5





- Resin model: hyperelastic-based formulation
 - Multiplicative decomposition $\mathbf{F} = \mathbf{F}^{ve} \cdot \mathbf{F}^{vp}, \quad \mathbf{C}^{ve} = \mathbf{F}^{ve^{T}} \cdot \mathbf{F}^{ve}, \quad J^{ve} = det(\mathbf{F}^{ve})$
 - Undamaged stress tensor definition
 - Elastic potential $\psi(\mathbf{C}^{\nu e})$
 - Undamaged first Piola-Kirchhoff stress tensor

$$\widehat{\mathbf{P}} = 2\mathbf{F}^{\mathbf{v}\mathbf{e}} \cdot \frac{\partial \psi(\mathbf{C}^{\mathbf{v}\mathbf{e}})}{\partial \mathbf{C}^{\mathbf{v}\mathbf{e}}} \cdot \mathbf{F}^{\mathbf{v}\mathbf{p}^{-T}}$$

- Undamaged Kirchhoff stress tensors
 - In current configuration

$$\widehat{\boldsymbol{\kappa}} = \widehat{\mathbf{P}} \cdot \mathbf{F}^T = 2\mathbf{F}^{\mathbf{v}e} \cdot \frac{\partial \psi(\mathbf{C}^{\nu e})}{\partial \mathbf{C}^{\mathbf{v}e}} \cdot \mathbf{F}^{\mathbf{v}e^T}$$

In co-rotational space

$$\widehat{\boldsymbol{\tau}} = \mathbf{C}^{\mathrm{ve}} \cdot \mathbf{F}^{\mathrm{ve}^{-1}} \widehat{\boldsymbol{\kappa}} \cdot \mathbf{F}^{\mathrm{ve}^{-T}} = 2\mathbf{C}^{\boldsymbol{\nu}\mathrm{e}} \cdot \frac{\partial \psi(\mathbf{C}^{\mathrm{ve}})}{\partial \mathbf{C}^{\mathrm{ve}}}$$

- Apparent stress tensor
 - Piola-Kirchhoff stress

$$\mathbf{P} = (\mathbf{1} - D_s) \big(\mathbf{1} - D_f \big) \widehat{\mathbf{P}}$$





- Resin model: logarithmic visco-elasticity
 - Elastic potentials ψ_i :

$$\psi_i(\mathbf{C}^{\mathrm{ve}}) = \frac{K_i}{2} \ln^2(J^{\mathrm{ve}}) + \frac{G_i}{4} (\ln(\mathbf{C}^{\mathrm{ve}}))^{\mathrm{dev}} : (\ln(\mathbf{C}^{\mathrm{ve}}))^{\mathrm{dev}}$$

- Dissipative potentials Υ_i

$$\Upsilon_i(\mathbf{C}^{\text{ve}}, \mathbf{q}_i) = -\mathbf{q}_i \colon \ln(\mathbf{C}^{\text{ve}}) + \left[\frac{1}{18K_i} \operatorname{tr}^2(\mathbf{q}_i) + \frac{1}{4G_i} \mathbf{q}_i^{\text{dev}} : \mathbf{q}_i^{\text{dev}}\right]$$

$$\begin{bmatrix} \dot{\mathbf{q}}_i^{\text{dev}} = \frac{2G_i}{g_i} & (\ln(\mathbf{C}^{\text{ve}}))^{\text{dev}} - \frac{1}{g_i} \mathbf{q}_i^{\text{dev}} \\ \text{tr} (\dot{\mathbf{q}}_i) = \frac{3K_i}{k_i} & \ln^2(J^{\text{ve}}) - \frac{1}{k_i} \text{tr} (\mathbf{q}_i) \end{bmatrix}$$



- Total potential ψ :

$$\begin{cases} \psi(\mathbf{C}^{ve}; \boldsymbol{q}_i) = \psi_{\infty}(\mathbf{C}^{ve}) + \sum_i [\psi_i(\mathbf{C}^{ve}) + \Upsilon_i(\mathbf{C}^{ve}, \mathbf{q}_i)] \\ \widehat{\mathbf{P}} = 2\mathbf{F}^{ve} \cdot \frac{\partial \psi(\mathbf{C}^{ve})}{\partial \mathbf{C}^{ve}} \cdot \mathbf{F}^{vp^{-T}} \end{cases}$$



- Resin model: visco-plasticity
 - Stress, back-stress $\boldsymbol{\varphi} = \hat{\boldsymbol{\tau}} - \hat{\boldsymbol{b}}$
 - Perzina plastic flow rule

 $\mathbf{D}^{\mathrm{vp}} = \dot{\mathbf{F}}^{\mathrm{vp}} \cdot \mathbf{F}^{\mathrm{vp}} = \frac{1}{n} \langle \phi \rangle^{\frac{1}{p}} \frac{\partial P}{\partial \hat{\tau}}$

Pressure dependent yield surface

$$\begin{bmatrix} \phi = \left(\frac{\varphi^{\text{eq}}}{\sigma_c}\right)^{\alpha} - \frac{m^{\alpha} - 1}{m+1} \frac{\text{tr}\boldsymbol{\varphi}}{\sigma_c} - \frac{m^{\alpha} + m}{m+1} \le 0\\ m = \frac{\sigma_t}{\sigma_c} \end{bmatrix}$$

Non-associated flow potential

$$P = (\varphi^{\rm eq})^2 + \beta \left(\frac{{\rm tr}\boldsymbol{\varphi}}{3}\right)^2$$

Equivalent plastic strain rate:

$$\dot{\boldsymbol{\gamma}} = \frac{\sqrt{\mathbf{D}^{\text{vp}}:\mathbf{D}^{\text{vp}}}}{\sqrt{\mathbf{1}+2\boldsymbol{v}_p^2}}$$
$$\boldsymbol{v}_p = \frac{9-2\beta}{18+2\beta}$$





- Resin model: failure softening - Failure surface $\begin{cases}
 \phi_f = \gamma - a \exp\left(-b \frac{\operatorname{tr}(\hat{r})}{3\hat{t}^{eq}}\right) - c \\
 \phi_f - r \le 0; \dot{r} \ge 0; \text{and } \dot{r}(\phi_f - r) = 0 \\
 \dot{\gamma}_f = \dot{r}
 \end{cases}$ - Damage evolution $\begin{cases}
 \dot{D}_f = H_f(\chi_f)^{\zeta_f} (1 - D_f)^{-\zeta_d} \dot{\chi}_f \\
 \chi_f = \max_{\tau} (\tilde{\gamma}_f(\tau)) \\
 \tilde{\gamma}_f - l_f^2 \Delta \tilde{\gamma}_f = \gamma_f \\
 l_f = 3 \mu m \quad V_0 \tilde{\gamma}_f \cdot \mathbf{N} = 0
 \end{cases}$
 - Affect ductility









- Composite model: Validation
 - Compression test



- PDR T.1015.14 project
 - ULiège, UCL (Belgium)
- Publications
 - <u>10.1016/j.ijsolstr.2016.06.008</u>
 - <u>10.1016/j.mechmat.2019.02.017</u>



Beginning

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Stochastic Homogenization of Composite Materials

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• Multi-scale modeling

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Material response Macro-scale BVP BVP

• For structures not several orders larger than the micro-structure size $L_{macro} >> L_{VE} >\sim L_{micro}$

For accuracy: Size of the mesoscale volume element smaller than the characteristic length of the macro-scale loading

Meso-scale volume element no longer statistically representative: • Stochastic Volume Elements



Beginning

• Material uncertainties affect structural behaviors





Beginning

Proposed methodology for material:

 To develop a stochastic Mean Field Homogenization method able to predict the probabilistic distribution of material response at an intermediate scale from microstructural constituents characterization





- Micro-structure stochastic model
 - 2000x and 3000x SEM images



Fibers detection







Micro-structure stochastic model

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Dependent variables generated using their empirical copula
 SEM sample
 Generated sample



Directly from copula generator

Statistic result from generated SVE

32

Beginning

0.8

Micro-structure stochastic model

- Dependent variables generated using their empirical copula
- Fiber additive process
 - 1) Define *N* seeds with first and second neighbors distances
 - 2) Generate first neighbor with its own first and second neighbors distances
 - 3) Generate second neighbor with its own first and second neighbors distances
 - 4) Change seeds & then change central fiber of the seeds



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- Micro-structure stochastic model
 - Arbitrary size
 - Arbitrary number







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Beginning

• Stochastic homogenization of SVEs

- Extraction of Stochastic Volume Elements
 - 2 sizes considered: $l_{\rm SVE} = 10 \ \mu m$ & $l_{\rm SVE} = 25 \ \mu m$
 - Window technique to capture correlation

$$R_{\mathbf{rs}}(\boldsymbol{\tau}) = \frac{\mathbb{E}\left[\left(r(\boldsymbol{x}) - \mathbb{E}(r)\right)\left(s(\boldsymbol{x} + \boldsymbol{\tau}) - \mathbb{E}(s)\right)\right]}{\sqrt{\mathbb{E}\left[\left(r - \mathbb{E}(r)\right)^{2}\right]}\sqrt{\mathbb{E}\left[\left(s - \mathbb{E}(s)\right)^{2}\right]}}$$

- For each SVE
 - Extract apparent homogenized material tensor \mathbb{C}_{M}

$$\begin{cases} \boldsymbol{\varepsilon}_{\mathrm{M}} = \frac{1}{V(\omega)} \int_{\omega} \boldsymbol{\varepsilon}_{\mathrm{m}} d\omega \\ \boldsymbol{\sigma}_{\mathrm{M}} = \frac{1}{V(\omega)} \int_{\omega} \boldsymbol{\sigma}_{\mathrm{m}} d\omega \\ \mathbb{C}_{\mathrm{M}} = \frac{\partial \boldsymbol{\sigma}_{\mathrm{M}}}{\partial \boldsymbol{u}_{\mathrm{M}} \otimes \boldsymbol{\nabla}_{\mathrm{M}}} \end{cases}$$

- Consistent boundary conditions:
 - Periodic (PBC)
 - Minimum kinematics (SUBC)
 - Kinematic (KUBC)





Beginning

• Stochastic homogenization of SVEs



Apparent properties

When l_{SVE} increases

- Average values for different BCs get closer (to PBC one)
- Distributions narrow
- Distributions get closer to normal



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• Stochastic homogenization of SVEs







Beginning

- Inverse stochastic identification
 - Comparison of homogenized properties from SVE realizations and stochastic MFH









- Non-linear inverse identification
 - Comparison SVE vs. MFH





- Damage-enhanced Mean-Field-homogenization
 - Virtual elastic unloading from previous state
 - Composite material unloaded to reach the stressfree state
 - Residual stress in components
 - Define Linear Comparison Composite
 - From elastic state

 $\Delta \boldsymbol{\epsilon}_{I/0}^{r} = \Delta \boldsymbol{\epsilon}_{I/0} + \Delta \boldsymbol{\epsilon}_{I/0}^{unload}$

Incremental-secant loading

$$\begin{cases} \boldsymbol{\sigma}_{\mathrm{M}} = \overline{\boldsymbol{\sigma}} = v_{0}\boldsymbol{\sigma}_{0} + v_{\mathrm{I}}\boldsymbol{\sigma}_{\mathrm{I}} \\ \boldsymbol{\Delta}\boldsymbol{\varepsilon}_{\mathrm{M}}^{\mathbf{r}} = \overline{\boldsymbol{\Delta}}\overline{\boldsymbol{\varepsilon}} = v_{0}\boldsymbol{\Delta}\boldsymbol{\varepsilon}_{0}^{\mathbf{r}} + v_{\mathrm{I}}\boldsymbol{\Delta}\boldsymbol{\varepsilon}_{\mathrm{I}}^{\mathbf{r}} \\ \boldsymbol{\Delta}\boldsymbol{\varepsilon}_{\mathrm{I}}^{\mathbf{r}} = \mathbb{B}^{\varepsilon} \big(\mathbf{I}, (1 - D_{0})\mathbb{C}_{0}^{\mathrm{S}}, \mathbb{C}_{\mathrm{I}}^{\mathrm{S}} \big) : \boldsymbol{\Delta}\boldsymbol{\varepsilon}_{0}^{\mathbf{r}} \end{cases}$$

Incremental secant operator

$$\Delta \boldsymbol{\sigma}_{\mathrm{M}} = \mathbb{C}_{\mathrm{M}}^{\mathrm{S}} \big(\mathrm{I}, (1 - D_0) \mathbb{C}_0^{\mathrm{S}}, \mathbb{C}_{\mathrm{I}}^{\mathrm{S}}, \boldsymbol{v}_{\mathrm{I}} \big) : \Delta \boldsymbol{\varepsilon}_{\mathrm{M}}^{\mathrm{r}}$$





- Damage-enhanced inverse identification
 - Comparison SVE vs. MFH





Generation of random field



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Beginning

• One single ply loading realization

- Random field and finite elements discretizations
- Non-uniform homogenized stress distributions
- Creates damage localization



• Ply loading realizations

- Simple failure criterion at (homogenized stress) loss of ellipticity
- Discrepancy in failure point





- STOMMMAC M.ERA-NET project (MFH for elasto-visco-plastic composites)
 - e-Xstream, ULiège (Belgium)
 - BATZ (Spain)
 - JKU, AC (Austria)
 - U Luxembourg (Luxemburg)
- Publications (doi)
 - <u>10.1016/j.compstruct.2018.01.051</u>
 - <u>10.1002/nme.5903</u>
 - <u>10.1016/j.cma.2019.01.016</u>



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Bayesian identification of stochastic Mean-Field Homogenization model parameters

STOMMMAC The research has been funded by the Walloon Region under the agreement no 1410246-STOMMMAC (CT-INT 2013-03-28) in the context of M-ERA.NET Joint Call 2014.



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• Multi-scale modeling

- 2 problems are solved concurrently
 - The macro-scale problem
 - The meso-scale problem (on a meso-scale
 - Volume Element)



Identification: Requires identification of micro-scale geometrical and material model parameters



Beginning

Proposed methodology

 To develop a stochastic Mean Field Homogenization method whose missing microconstituents properties are inferred from coupons tests





Beginning

- Fibre distribution effect
 - 2-step homogenization



- For uniaxial tests along direction θ : $\sigma_M = \sigma_M (I(\psi(p)), \mathbb{C}_0, \mathbb{C}_I; \theta, \varepsilon_M)$





- Fibre distribution effect
 - Skin-core effect





Beginning

• Experimental characterization Fiber orientation and aspect ratio (JKU)



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Composite material response (BATZ)



- Assume a distribution of the matrix Young's modulus
 - Beta distribution $E_0 \sim \beta_{\alpha,\beta,a,b}$ with $\beta_{\alpha,\beta,a,b}(y) = \frac{(y-a)^{\alpha-1}(y-b)^{\beta-1}}{(b-a)^{\alpha+\beta+1}B(\alpha,\beta)}$
 - Matrix Young 's modulus corresponding to experimental measurements
 - $E_{0c}^{(n)}$ with $n = 1..n_{\text{total}}$, for all directions and positions
 - Bayes' theorem

 $\pi_{\text{post}}(\alpha, \beta, a, b | \hat{E}_{0c}) \propto \pi(\hat{E}_{0c} | \alpha, \beta, a, b) \pi_{\text{prior}}(\alpha) \pi_{\text{prior}}(\beta) \pi_{\text{prior}}(a) \pi_{\text{prior}}(b)$ • Priors: $\pi_{\text{prior}}(x) = \Gamma_{\alpha,\beta,a,c}$ with $\Gamma_{\alpha,\beta,a,c}(y) = \frac{\left(\frac{y-a}{c}\right)^{\alpha-1} \beta^{\alpha} e^{-\beta\left(\frac{y-a}{c}\right)}}{c\Gamma(\alpha)}$

• Likelihood:
$$\pi(\hat{E}_{0c}|\alpha,\beta,a,b) = \prod_{n=1}^{n_{\text{total}}} \beta_{\alpha,\beta,a,b}(E_{0c}^{(n)})$$

$$\prod_{n=1}^{n_{\text{total}}} \beta_{\alpha,\beta,a,c} \left(E_{0c}^{(n)} \right) \pi_{\text{prior}}(\alpha) \pi_{\text{prior}}(\beta) \pi_{\text{prior}}(a) \pi_{\text{prior}}(b)$$



Beginning

• Assume a distribution of the matrix Young's modulus

- Inference:
$$\pi_{\text{post}}(\alpha, \beta, a, b | \hat{E}_{0c}) \propto \prod_{n=1}^{n_{\text{total}}} \beta_{\alpha, \beta, a, c} \left(E_{0c}^{(n)} \right) \pi_{\text{prior}}(\alpha) \pi_{\text{prior}}(\beta) \pi_{\text{prior}}(a) \pi_{\text{prior}}(b)$$

• $i = 1..n_{pos}$, with n_{pos} the number of positions tested (5, positions #1-#5)





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•(5)

6

•3

•(4)

•2

1

Validation

- Evaluate stochastic response at Position 6
 - Perform stochastic homogenization from $\pi_{\text{post}}(\alpha, \beta, a, b | \hat{E}_{0c})$
 - From sampling of $[\alpha, \beta, a, b]$, evaluate $E_0 \sim \beta_{\alpha, \beta, a, b}$
 - From sampling of $[E_0]$, evaluate composite response

 $E_{\rm MFH} = E_{\rm MFH} \big(I(\psi(\boldsymbol{p}), a_r), E_0 , \mathbb{C}_{\rm I} ; \boldsymbol{\theta} \big)$

• Compare with experimental measurements $\hat{E}_c^{(6,j)}$



- Extension to non-linear behavior
 - More parameters to infer
 - Matrix Young's modulus E_0
 - Matrix yield stress σ_{Y_0}
 - Matrix hardening law $R(p_0) = h p_0^{m_1} (1 - \exp(-m_2 p_0))$
 - Effective aspect ratio a_r
 - 2-Step MFH model requires many iterations
 - Incremental secant approach

$$\begin{cases} \boldsymbol{\sigma}_{\mathrm{M}} = \overline{\boldsymbol{\sigma}} = v_{0}\boldsymbol{\sigma}_{0} + v_{\mathrm{I}}\boldsymbol{\sigma}_{\mathrm{I}} \\ \boldsymbol{\Delta}\boldsymbol{\varepsilon}_{\mathrm{M}}^{\mathbf{r}} = \overline{\boldsymbol{\Delta}\boldsymbol{\varepsilon}} = v_{0}\boldsymbol{\Delta}\boldsymbol{\varepsilon}_{0}^{\mathbf{r}} + v_{\mathrm{I}}\boldsymbol{\Delta}\boldsymbol{\varepsilon}_{\mathrm{I}}^{\mathbf{r}} \\ \boldsymbol{\Delta}\boldsymbol{\varepsilon}_{\mathrm{I}}^{\mathbf{r}} = \mathbb{B}^{\varepsilon}(\mathrm{I}, \mathbb{C}_{0}^{\mathrm{S}}, \mathbb{C}_{\mathrm{I}}^{\mathrm{S}}): \boldsymbol{\Delta}\boldsymbol{\varepsilon}_{0}^{\mathbf{r}} \end{cases}$$

Too expensive for BI

Definition of parameters





- Speed up the evaluation of the likelihood
 - Likelihood
 - $\boldsymbol{\pi}(\hat{\sigma}_{\mathrm{M}}(t)|[\varepsilon_{\mathrm{M}}(t'\leq t),\boldsymbol{\vartheta}])$
 - With $\boldsymbol{\vartheta} = [E_0, \sigma_{Y_0}, h, m_1, m_2, a_r]$
 - 2-Step MFH model $\sigma_{\rm MFH}(t)$ $= \sigma_{\rm MFH} (I(\psi(\mathbf{p}), a_r), E_0)$



- Too expensive for BI
- Use of a surrogate
 - $\sigma_{\text{NNW}}(t) = \sigma_{\text{NNW}}(\boldsymbol{\varepsilon}_{\mathbf{M}}(t), \boldsymbol{\vartheta}, \mathbb{C}_{\mathbf{I}}; \boldsymbol{\theta})$
 - Constructed using artificial Neural Network
 - Trained fusing the 2-Step MFH model

 $\begin{aligned} \sigma_{\rm MFH}(t) \\ &= \sigma_{\rm MFH} \big({\rm I}(\psi(\boldsymbol{p}), a_r), E_0 \\ , \qquad & \mathbb{C}_{\rm I} \ , \varepsilon_{\rm M}(t') \end{aligned}$







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- Assume a noise in the measurements & use surrogate model
 - Measurements at strain *i* in direction θ_i :

$$\Sigma_{c}^{(i,j,k)} = \sigma_{\text{NNW}}^{(i,j)} \left(\boldsymbol{\varepsilon}_{\mathbf{M}}^{(i,j)}, \boldsymbol{\vartheta}, \mathbb{C}_{\mathbf{I}}; \boldsymbol{\theta}_{j} \right) + \text{noise}^{(i,j)}$$

$$\pi \left(\Sigma_{c}^{(i,j,k)} | \left[\boldsymbol{\varepsilon}_{\mathbf{M}}^{(i,j)}, \boldsymbol{\vartheta} \right] \right)$$

= $\pi_{\text{noise}}^{(i,j)} \left(\Sigma_{c}^{(i,j,k)} - \sigma_{\text{NNW}}^{(i,j)} \left(\boldsymbol{\varepsilon}_{\mathbf{M}}^{(i,j)}, \boldsymbol{\vartheta}, \mathbb{C}_{\mathrm{I}}; \boldsymbol{\theta}_{j} \right) \right)$

• $j = 1..n_{dir}$, with

 $n_{\rm dir}$ the number of directions θ_i tested

•
$$i = 1..n_{\varepsilon}^{(j)}$$
, with

 n_{ε} the number of stress-strain points

•
$$k = 1..n_{\text{test}}^{(i,j)}$$
, with

 $n_{\text{test}}^{(i,j)}$ the number of samples tested at point *i* along direction θ_j

- Noise function from $n_{\text{test},i,j}$ measurements at strain *i* in direction θ_j :

$$\pi_{\text{noise}^{(i,j)}}(y) = \frac{1}{\sqrt{2\pi}} \sigma_{\Sigma_c^{(i,j)}} \exp\left(-\frac{y^2}{2\sigma_{\Sigma_c^{(i,j)}}^2}\right)$$

Bayes' theory:

 $\pi_{\text{post}}(\boldsymbol{\vartheta}|\boldsymbol{\hat{\varepsilon}}_{M},\boldsymbol{\hat{\Sigma}}_{c}) \propto \pi_{\text{prior}}(\boldsymbol{\vartheta}) \prod_{j=1}^{n_{\text{dir}}} \prod_{i=1}^{n_{\varepsilon}^{(i)}} \prod_{k=1}^{n_{\text{test}}^{(i,j)}} \pi_{\text{noise}}^{(i,j)} \left(\Sigma_{c}^{(i,j,k)} - \sigma_{\text{NNW}}^{(i,j)} \left(\boldsymbol{\varepsilon}_{M}^{(i,j)}, \boldsymbol{\vartheta}, \mathbb{C}_{I}; \boldsymbol{\theta}_{j} \right) \right)$





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Results

 $\pi_{\text{post}}(\boldsymbol{\vartheta}|\boldsymbol{\hat{\varepsilon}}_{M},\boldsymbol{\hat{\Sigma}}_{C}) \propto \pi_{\text{prior}}(\boldsymbol{\vartheta}) \quad \prod_{j=1}^{n_{\text{dir}}} \prod_{i=1}^{n_{\varepsilon}^{(j)}} \prod_{k=1}^{n_{\text{test}}^{(i,j)}} \pi_{\text{noise}}^{(i,j)} \left(\boldsymbol{\Sigma}_{c}^{(i,j,k)} - \sigma_{\text{NNW}}^{(i,j)} \left(\boldsymbol{\varepsilon}_{M}^{(i,j)}, \boldsymbol{\vartheta}, \boldsymbol{\mathbb{C}}_{I}; \boldsymbol{\theta}_{j} \right) \right)$





Beginning

• Verification

$$\pi_{\text{post}}(\boldsymbol{\vartheta}|\boldsymbol{\hat{\varepsilon}}_{M},\boldsymbol{\hat{\Sigma}}_{C}) \propto \pi_{\text{prior}}(\boldsymbol{\vartheta}) \prod_{j=1}^{n_{\text{dir}}} \prod_{i=1}^{n_{\varepsilon}^{(j)}} \prod_{k=1}^{n_{\text{test}}^{(i,j)}} \pi_{\text{noise}}^{(i,j)} \left(\Sigma_{c}^{(i,j,k)} - \sigma_{\text{NNW}}^{(i,j)} \left(\boldsymbol{\varepsilon}_{M}^{(i,j)}, \boldsymbol{\vartheta}, \mathbb{C}_{I}; \boldsymbol{\theta}_{j} \right) \right)$$





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- Publications (doi)
 - <u>10.1016/j.compstruct.2019.03.066</u>





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Non-Local Damage Mean-Field-Homogenization

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 $L_{\text{macro}} >> L_{\text{VE}} >> L_{\text{micro}}$

For accuracy: Size of the mesoscale volume element smaller than the characteristic length of the macro-scale loading To be statistically representative: Size of the meso-scale volume element larger than the characteristic length of the microstructure



- Materials with strain softening
 - Incremental forms
 - Strain increments in the same direction

 $\Delta \boldsymbol{\varepsilon}_{\mathrm{I}} = \mathbf{B}^{\varepsilon} \left(\mathrm{I}, \, \bar{\mathbf{C}}_{0}^{\mathrm{alg}}, \, \bar{\mathbf{C}}_{\mathrm{I}}^{\mathrm{alg}} \right) : \Delta \boldsymbol{\varepsilon}_{0}$

 Because of the damaging process, the fiber phase is elastically unloaded during matrix softening

- Solution: new incremental-secant method
 - We need to define the LCC from another stress state





- Based on the incremental-secant approach
 - Perform a virtual elastic unloading from previous solution
 - Composite material unloaded to reach the stress-free state
 - Residual stress in components
 - Apply MFH from unloaded state
 - New strain increments (>0)

 $\Delta \pmb{\epsilon}_{I/0}^r = \Delta \pmb{\epsilon}_{I/0} + \Delta \pmb{\epsilon}_{I/0}^{unload}$

Use of secant operators

$$\Delta \boldsymbol{\varepsilon}_{\mathrm{I}}^{\mathrm{r}} = \mathbf{B}^{\varepsilon} \big(\mathrm{I}, (1-D) \bar{\mathbf{C}}_{0}^{\mathrm{Sr}}, \bar{\mathbf{C}}_{\mathrm{I}}^{\mathrm{S0}} \big) : \Delta \boldsymbol{\varepsilon}_{0}^{\mathrm{r}}$$

Possibility of unloading

$$\begin{cases} \Delta \boldsymbol{\epsilon}_{\mathrm{I}}^{\mathrm{r}} > \boldsymbol{0} \\ \Delta \boldsymbol{\epsilon}_{\mathrm{I}} < \boldsymbol{0} \end{cases}$$





Beginning

- New results for damage
 - Fictitious composite
 - 50%-UD fibres
 - Elasto-plastic matrix with damage









Beginning

- Material models
 - Elasto-plastic material
 - Stress tensor $\boldsymbol{\sigma} = \boldsymbol{C}^{el}: (\boldsymbol{\epsilon} \boldsymbol{\epsilon}^{pl})$
 - Yield surface $f(\boldsymbol{\sigma}, p) = \boldsymbol{\sigma}^{eq} \sigma^Y R(p) \le 0$
 - Plastic flow $\Delta \epsilon^{\mathbf{pl}} = \Delta p \mathbf{N}$ & $\mathbf{N} = \frac{\partial f}{\partial \sigma}$
 - Local damage model
 - Apparent-effective stress tensors $\boldsymbol{\sigma} = (1 D) \widehat{\boldsymbol{\sigma}}$
 - Plastic flow in the effective stress space
 - Damage evolution $\Delta D = F_D(\mathbf{\epsilon}, \Delta p)$
 - Non-Local damage model [Peerlings et al., 1996]
 - Damage evolution $\Delta D = F_D(\mathbf{\epsilon}, \Delta \tilde{p})$
 - Anisotropic governing equation $\tilde{p} \nabla \cdot (\mathbf{c}_{\mathbf{g}} \cdot \nabla \tilde{p}) = p$





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Beginning

Laminate studies

- Bulk material law
 - Non-local damage-enhanced MFH
 - Intra-laminar failure
 - Account for anisotropy
- Interface
 - DG/Cohesive zone model
 - Inter-laminar failure





Beginning

• $[45^{\circ}_4 / -45^{\circ}_4]_{s}$ - open hole laminate (epoxy- with 60% UD CF)



• $[90^{\circ}/45^{\circ}/-45^{\circ}/90^{\circ}/0^{\circ}]_{s}$ - open hole laminate

Intra-laminar failure along fiber directions (experiments: IMDEA Materials)


Non-Local Damage Mean-Field-Homogenization

- $[90^{\circ}/45^{\circ}/-45^{\circ}/90^{\circ}/0^{\circ}]_{s}$ open hole laminate
 - Inter-laminar failure compared to experimental results (experiments: IMDEA Materials)



Non-Local Damage Mean-Field-Homogenization

- SIMUCOMP ERA-NET project
 - e-Xstream, CENAERO, ULiège (Belgium)
 - IMDEA Materials (Spain)
 - CRP Henri-Tudor (Luxemburg)
- Publications (doi)
 - <u>10.1016/j.compstruct.2015.02.070</u>
 - <u>10.1016/j.ijsolstr.2013.07.022</u>
 - <u>10.1016/j.ijplas.2013.06.006</u>
 - <u>10.1016/j.cma.2012.04.011</u>
 - <u>10.1007/978-1-4614-4553-1_13</u>



Computational & Multiscale Mechanics of Materials



Boundary conditions and tangent operator in multiphysics computational homogenization

ARC 09/14-02 BRIDGING - From imaging to geometrical modelling of complex micro structured materials: Bridging computational engineering and material science

The authors gratefully acknowledge the financial support from F.R.S-F.N.R.S. under the project number PDR T.1015.14



April 2019 CM3 research projects

- Multi-scale modeling
 - 2 problems are solved concurrently
 - The macro-scale problem
 - The meso-scale problem (on a meso-scale Volume Element)







For accuracy: Size of the mesoscale volume element smaller than the characteristic length of the macro-scale loading To be statistically representative: Size of the meso-scale volume element larger than the characteristic length of the microstructure



April 2019 CM3 research projects

- Generalized multi-physics representation
 - Strong form $\mathcal{P} \cdot \mathcal{V}_0 = 0$
 - Fully-coupled constitutive law $\mathcal{P} = \mathcal{P}(\mathcal{X}^{C}, \mathcal{F}; \mathcal{Z})$
 - \mathcal{F} : generalized deformation gradient, $\mathcal{X}^{\mathcal{C}}$: fields appearing in the constitutive relations
 - Z: internal variables

• Tangent operators $\mathcal{L} = \frac{\partial \mathcal{P}}{\partial \mathcal{F}} \& \mathcal{J} = \frac{\partial \mathcal{P}}{\partial \chi^c}$ but also $\mathcal{Y}_{\mathcal{F}} = \frac{\partial Z}{\partial \mathcal{F}} \& \mathcal{Y}_{\chi^c} = \frac{\partial Z}{\partial \chi^c}$





- Generalized microscopic boundary conditions
 - Arbitrary field k kinematics: $\mathcal{X}_{m}^{k} = \mathcal{X}_{M}^{k} + \mathcal{F}_{M}^{k} \cdot X_{m} + \mathcal{W}_{m}^{k}$
 - Constrained field k equivalence:

e:
$$\int_{\omega_0} C_m^k \mathcal{X}_m^{C^k} d\omega = \int_{\omega_0} C_m^k d\omega \mathcal{X}_M^{C^k}$$

- E.g. periodic boundary conditions

Define an interpolant map

$$\mathbb{S}^i = \sum \mathbb{N}^i_k(\boldsymbol{X}_m) a^i_k$$

Substitute fluctuation fields

$$W_m^k(X_m^+) = \mathbb{S}^i(X_m^-) = W_m^k(X_m^-)$$



Boundary nodeControl node

Fluctuation





Beginning

Microscale BVP

Weak formulation

$$\begin{cases} \mathcal{P}_{\mathrm{m}} \cdot \nabla_{0} = 0 & \text{with } \mathcal{P}_{m}(\mathcal{X}_{m}^{C}, \mathcal{F}_{m}; \mathcal{Z}_{m} \\ \mathcal{X}_{\mathrm{m}}^{k} = \mathcal{X}_{\mathrm{M}}^{k} + \mathcal{F}_{\mathrm{M}}^{k} \cdot \mathcal{X}_{\mathrm{m}} + \mathcal{W}_{\mathrm{m}}^{k} \\ \int_{\omega_{0}} \mathcal{C}_{m}^{k} \mathcal{X}_{\mathrm{m}}^{C^{k}} d\omega = \int_{\omega_{0}} \mathcal{C}_{m}^{k} d\omega \, \mathcal{X}_{\mathrm{M}}^{C^{k}} \end{cases} \end{cases}$$

- Weak finite element constrained form $(\omega_0 = \cup_e \omega^e)$

$$\begin{cases} \mathbf{f}_{\mathrm{m}}(\boldsymbol{\mathcal{U}}_{m}) - \mathbf{C}^{\mathrm{T}}\boldsymbol{\lambda} = 0\\ \mathbf{C}\boldsymbol{\mathcal{U}}_{m} - \mathbf{S}\begin{bmatrix} \boldsymbol{\mathcal{F}}_{\mathrm{M}}\\ \boldsymbol{\mathcal{X}}_{\mathrm{M}}^{C} \end{bmatrix} = 0 \end{cases}$$

System linearization

$$\mathbf{Q}^{\mathrm{T}} \frac{\partial \mathbf{f}_{\mathrm{m}}}{\partial \boldsymbol{\mathcal{U}}_{m}} \mathbf{Q} \delta \boldsymbol{\mathcal{U}}_{m} + \mathbf{r} - \mathbf{Q}^{\mathrm{T}} \frac{\partial \mathbf{f}_{\mathrm{m}}}{\partial \boldsymbol{\mathcal{U}}_{m}} \mathbf{C}^{\mathrm{T}} (\mathbf{C}\mathbf{C}^{\mathrm{T}})^{-1} \left(\mathbf{r}_{c} - \mathbf{S} \begin{bmatrix} \delta \mathcal{F}_{\mathrm{M}} \\ \delta \boldsymbol{\chi}_{\mathrm{M}}^{C} \end{bmatrix} \right) = 0$$
$$\mathbf{C} \delta \boldsymbol{\mathcal{U}}_{m} + \mathbf{r}_{\mathrm{c}} - \mathbf{S} \begin{bmatrix} \delta \mathcal{F}_{\mathrm{M}} \\ \delta \boldsymbol{\chi}_{\mathrm{M}}^{C} \end{bmatrix} = 0 \qquad \& \qquad \mathbf{Q} = \mathbf{I} - \mathbf{C}^{\mathrm{T}} (\mathbf{C}\mathbf{C}^{\mathrm{T}})^{-1} \mathbf{C}$$



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- Multi-scale resolution
 - System linearization

$$\begin{cases} \mathbf{Q}^{\mathrm{T}} \frac{\partial \mathbf{f}_{\mathrm{m}}}{\partial \boldsymbol{\mathcal{U}}_{m}} \mathbf{Q} \delta \boldsymbol{\mathcal{U}}_{m} + \mathbf{r} - \mathbf{Q}^{\mathrm{T}} \frac{\partial \mathbf{f}_{\mathrm{m}}}{\partial \boldsymbol{\mathcal{U}}_{m}} \mathbf{C}^{\mathrm{T}} (\mathbf{C} \mathbf{C}^{\mathrm{T}})^{-1} \left(\mathbf{r}_{c} - \mathbf{S} \begin{bmatrix} \delta \mathcal{F}_{\mathrm{M}} \\ \delta \boldsymbol{\mathcal{X}}_{\mathrm{M}}^{C} \end{bmatrix} \right) = 0 \\ \mathbf{C} \delta \boldsymbol{\mathcal{U}}_{m} + \mathbf{r}_{\mathrm{c}} - \mathbf{S} \begin{bmatrix} \delta \mathcal{F}_{\mathrm{M}} \\ \delta \boldsymbol{\mathcal{X}}_{\mathrm{M}}^{C} \end{bmatrix} = 0 \qquad \& \qquad \mathbf{Q} = \mathbf{I} - \mathbf{C}^{\mathrm{T}} (\mathbf{C} \mathbf{C}^{\mathrm{T}})^{-1} \mathbf{C} \end{cases}$$

- FEM resolution:
$$\delta \mathcal{F}_{M} = \delta \mathcal{X}_{M}^{C} = 0$$

 $\delta \mathcal{U}_{m} = -\tilde{K}^{-1} \left(\mathbf{r} + \left(\mathbf{C}^{T} - \mathbf{Q}^{T} \frac{\partial \mathbf{f}_{m}}{\partial \mathcal{U}_{m}} \mathbf{C}^{T} (\mathbf{C}\mathbf{C}^{T})^{-1} \right) \mathbf{r}_{c} \right)$
- Constraints effect: $\mathbf{r} = \mathbf{r}_{c} = 0$
 $\frac{\partial \mathcal{U}_{m}}{\partial \left[\mathcal{F}_{M} \quad \mathcal{X}_{M}^{C} \right]^{T}} = \tilde{K}^{-1} \left(\mathbf{C}^{T} - \mathbf{Q}^{T} \frac{\partial \mathbf{f}_{m}}{\partial \mathcal{U}_{m}} \mathbf{C}^{T} (\mathbf{C}\mathbf{C}^{T})^{-1} \right) \mathbf{S}$
- Only one matrix to factorize
 $\tilde{K} = \mathbf{C}^{T} \mathbf{C} + \mathbf{Q}^{T} \frac{\partial \mathbf{f}_{m}}{\partial \mathcal{U}_{m}} \mathbf{Q}$

Macro-scale operators at low cost

$$\begin{bmatrix} \frac{\partial \mathcal{P}_{M}}{\partial \mathcal{F}_{M}} & \frac{\partial \mathcal{P}_{M}}{\partial \mathcal{X}_{M}^{C}} \\ \frac{\partial \mathcal{Z}_{M}}{\partial \mathcal{F}_{M}} & \frac{\partial \mathcal{Z}_{M}}{\partial \mathcal{X}_{M}^{C}} \end{bmatrix} = \left(\bigwedge_{\omega^{e}} \frac{1}{V(\omega_{0})} \int_{\omega^{e}_{0}} \begin{bmatrix} \frac{\partial \mathcal{P}_{m}}{\partial \mathcal{F}_{m}} \mathbf{B}^{e} & \frac{\partial \mathcal{P}_{m}}{\partial \mathcal{X}_{m}^{C}} \mathbf{N}^{e} \\ \frac{\partial \mathcal{Z}_{m}}{\partial \mathcal{F}_{m}} \mathbf{B}^{e} & \frac{\partial \mathcal{Z}_{m}}{\partial \mathcal{X}_{m}^{C}} \mathbf{N}^{e} \end{bmatrix} d\omega \right) \frac{\partial \boldsymbol{u}_{m}}{\partial [\mathcal{F}_{M} & \mathcal{X}_{M}^{C}]^{T}}$$



Beginning







Beginning



- BRIDGING ARC project (Periodic boundary conditions)
 - ULiège, Applied Sciences (A&M, EEI, ICD)
 - ULiège, Sciences (CERM)
- PDR T.1015.14 project (MFH with second-order moments)
 - ULiège, UCL (Belgium)
- Publications
 - <u>10.1007/s00466-016-1358-z</u>
 - <u>10.1016/j.commatsci.2011.10.017</u>



Computational & Multiscale Mechanics of Materials





Computational Homogenization For Cellular Materials

ARC 09/14-02 BRIDGING - From imaging to geometrical modelling of complex micro structured materials: Bridging computational engineering and material science



April 2019 CM3 research projects

• Multi-scale modeling

- 2 problems are solved concurrently
 - The macro-scale problem
 - The meso-scale problem (on a meso-scale Volume Element)



- What if homogenized properties loose ellipticity?
 - Buckling of honeycomb structures





April 2019 CM3 research projects

• DG-based second-order FE²

- Macro-scale
 - High-order Strain-Gradient formulation
 - C¹ weakly enforced by DG
 - Partitioned mesh (//)
- Transition
 - Gauss points on different processors
 - Each Gauss point is associated to one mesh and one solver

- Micro-scale
 - Usual 3D finite elements
 - High-order periodic boundary conditions
 - Non-conforming mesh
 - Use of interpolant functions





Instabilities

- Micro-scale: buckling
- Macro-scale: localization bands
- Captured owing to
 - Second-order homogenization
 - Ad-hoc periodic boundary conditions
 - Path following method









• Open-hole plate





BRIDGING ARC project

- ULiège, Applied Sciences (A&M, EEI, ICD)
- ULiège, Sciences (CERM)
- Publications
 - <u>10.1016/j.mechmat.2015.07.004</u>
 - <u>10.1016/j.ijsolstr.2014.02.029</u>
 - <u>10.1016/j.cma.2013.03.024</u>



Computational & Multiscale Mechanics of Materials





Stochastic 3-Scale Models for Polycrystalline Materials

3SMVIB: The research has been funded by the Walloon Region under the agreement no 1117477 (CT-INT 2011-11-14) in the context of the ERA-NET MNT framework.



April 2019 CM3 research projects

• Multi-scale modeling

- 2 problems are solved concurrently
 - The macro-scale problem
 - The meso-scale problem (on a meso-scale Volume Element)



• For structures not several orders larger than the micro-structure size $L_{macro} >> L_{VE} >\sim L_{micro}$

For accuracy: Size of the mesoscale volume element smaller than the characteristic length of the macro-scale loading

Meso-scale volume element no longer statistically representative: • Stochastic Volume Elements



Beginning

• Key idea

Micro-scale	Meso-scale	Macro-scale	
 Samples of stochastic volume elements Random microstructure 	 > Intermediate scale > The distribution of the material property ℙ(C) is defined 	 Uncertainty quantification of the macro-scale quantity Quantity of interest distribution P(0) 	
Stochastic Contraction	n Mean value of material property SVE size Variance of material property SVE size	Probability density Quantity of interest	
université	April 2019 CM3 research projects	91 <u>Beginning</u>	

- Material structure: grain orientation distribution
 - Grain orientation by XRD (X-ray Diffraction) measurements on 2 µm-thick poly-silicon films



XRD images provided by IMT Bucharest, Rodica Voicu, Angela Baracu, Raluca Muller



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- Application to polycrystalline materials: The micro-scale to meso-scale transition
 - Stochastic homogenization



$$\sigma_{m^{i}} = \mathbb{C}_{i}: \epsilon_{m^{i}} , \forall i$$
Stochastic
Homogenization
$$\sigma_{M} = \mathbb{C}_{M}: \epsilon_{M}$$
Samples of the meso-scale
homogenized elasticity tensors

- Homogenized Young's modulus distribution



- Application to polycrystalline materials: The meso-scale spatial correlation
 - Use of the window technique

$$R_{\mathbb{C}}^{(rs)}(\boldsymbol{\tau}) = \frac{\mathbb{E}\left[\left(\mathbb{C}^{(r)}(\boldsymbol{x}) - \mathbb{E}(\mathbb{C}^{(r)})\right)\left(\mathbb{C}^{(s)}(\boldsymbol{x}+\boldsymbol{\tau}) - \mathbb{E}(\mathbb{C}^{(s)})\right)\right]}{\sqrt{\mathbb{E}\left[\left(\mathbb{C}^{(r)} - \mathbb{E}(\mathbb{C}^{(r)})\right)^{2}\right]\mathbb{E}\left[\left(\mathbb{C}^{(s)} - \mathbb{E}(\mathbb{C}^{(s)})\right)^{2}\right]}}$$



- Definition of the correlation length



$$L_{\mathbb{C}}^{(rs)} = \frac{\int_{-\infty}^{\infty} R_{\mathbb{C}}^{(rs)}}{R_{\mathbb{C}}^{(rs)}(0)}$$



Beginning

- Application to polycrystalline materials: The meso-scale random field
 - Accounts for the meso-scale distribution & spatial correlation



Needs to be generated using a stochastic model

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- Stochastic model of Gaussian meso-scale random fields
 - Define the homogenous zero-mean random field $\mathcal{A}'(x, \theta)$
 - Elasticity tensor $\mathbb{C}_{M}(x,\theta)$ (matrix form C_{M}) is bounded
 - $\boldsymbol{\varepsilon}: (\mathbb{C}_{M} \mathbb{C}_{L}): \boldsymbol{\varepsilon} > 0 \qquad \forall \boldsymbol{\varepsilon}$
 - Use a Cholesky decomposition

$$\boldsymbol{C}_{\mathrm{M}}(\boldsymbol{x},\boldsymbol{\theta}) = \boldsymbol{C}_{\mathrm{L}} + \left(\bar{\boldsymbol{\mathcal{A}}} + \boldsymbol{\mathcal{A}}'(\boldsymbol{x},\boldsymbol{\theta})\right)^{\mathrm{T}} \left(\bar{\boldsymbol{\mathcal{A}}} + \boldsymbol{\mathcal{A}}'(\boldsymbol{x},\boldsymbol{\theta})\right)$$

Evaluate the covariance function

 $\tilde{R}_{\mathcal{A}'}^{(rs)}(\boldsymbol{\tau}) = \sigma_{\mathcal{A}'^{(r)}} \sigma_{\mathcal{A}'^{(s)}} R_{\mathcal{A}'}^{(rs)}(\boldsymbol{\tau})$ $= \mathbb{E}\left[\left(\mathcal{A}'^{(r)}(\boldsymbol{x}) \right) \left(\mathcal{A}'^{(s)}(\boldsymbol{x} + \boldsymbol{\tau}) \right) \right]$



- Evaluate the spectral density matrix from periodized zero-padded matrix $\widetilde{R}_{\mathcal{V}'}^{\mathrm{P}}(\tau)$ $S_{\mathcal{A}'}^{(rs)}[\omega^{(m)}] = \sum_{n} \widetilde{R}_{\mathcal{A}'}^{\mathrm{P}}{}^{(rs)}[\tau^{(n)}]e^{-2\pi i \tau^{(n)} \cdot \omega^{(m)}} \& S_{\mathcal{A}'}[\omega^{(m)}] = H_{\mathcal{A}'}[\omega^{(m)}]H_{\mathcal{A}'}^{*}[\omega^{(m)}]$
- Generate a Gaussian random field $\mathcal{A}'(x, \theta)$

$$\mathcal{A}^{\prime(r)}(\boldsymbol{x},\boldsymbol{\theta}) = \sqrt{2\Delta\omega} \,\Re\left(\sum_{s} \sum_{m} H_{\mathcal{A}^{\prime}}^{(rs)} [\boldsymbol{\omega}^{(m)}] \,\eta^{(s,m)} \,e^{2\pi i \left(\boldsymbol{x}\cdot\boldsymbol{\omega}^{(m)} + \boldsymbol{\theta}^{(s,m)}\right)}\right)$$



Beginning

- Stochastic model of non-Gaussian meso-scale random fields
 - Start from micro-sampling of the stochastic homogenization
 - The continuous form of the targeted PSD function

$$\boldsymbol{S}^{\mathrm{T}(rs)}(\boldsymbol{\omega}) = \boldsymbol{\Delta}\boldsymbol{\tau}\boldsymbol{S}^{(rs)}_{\boldsymbol{\mathcal{V}}'}[\boldsymbol{\omega}^{(m)}] = \boldsymbol{\Delta}\boldsymbol{\tau}\sum_{n} \widetilde{\boldsymbol{R}}^{\mathrm{P}}_{\boldsymbol{\mathcal{A}}'}{}^{(rs)}[\boldsymbol{\tau}^{(n)}]e^{-2\pi i\boldsymbol{\tau}^{(n)}\cdot\boldsymbol{\omega}^{(m)}}$$

- The targeted marginal distribution density function $F^{NG(r)}$ of the random variable $\mathcal{A}'^{(r)}$
- A marginal Gaussian distribution $F^{G(r)}$ of zero-mean and targeted variance $\sigma_{\mathcal{A}'^{(r)}}$
- Iterate





Beginning

- The meso-scale stochastic model
 - Application to film deposited at 610 °C:
 - Comparison between micro-samples and generated fields







- Application to polycrystalline materials: The meso-scale to macro-scale transition
 - Convergence in terms of $\alpha = \frac{l_{\mathbb{C}}}{l_{\text{mesh}}}$, the correlation length and macro-mesh ratio
 - The results converge

versité

- With the mesh size for all the SVE sizes
- Toward the direct Monte Carlo simulations results



- Application to polycrystalline materials: The meso-scale to macro-scale transition
 - Comparison with direct Monte Carlo simulations



Relative difference in the mean: 0.57 %





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Relative difference in the mean: 0.44%

Thermo-mechanical homogenization ۲ Х Down-scaling $\boldsymbol{\sigma}_{M}, \boldsymbol{q}_{M}, (\rho_{M}C_{\nu M})$ $\boldsymbol{\varepsilon}_{\mathrm{M}},$ $\mathbf{\varepsilon}_{\mathrm{M}} = \frac{1}{V(\omega)} \int_{\omega} \mathbf{\varepsilon}_{\mathrm{m}} d\omega$ $\mathbb{C}_{M}, \kappa_{M}, \boldsymbol{\alpha}_{M} \mathbb{C}_{M}, \boldsymbol{\alpha}_{M} \mathbb{C}$ $\nabla_{\mathrm{M}} \vartheta_{\mathrm{M}},$ ϑ_{M} $\nabla_{\rm M}\vartheta_{\rm M} = \frac{1}{V(\omega)} \int_{\omega} \nabla_{\rm m}\vartheta_{\rm m} d\omega$ Meso-scale BVP $\vartheta_{\rm M} = \frac{1}{V(\omega)} \int_{\omega} \frac{\rho_{\rm m} C_{\nu \rm m}}{\rho_{\rm M} C_{\nu \rm m}} \vartheta_{\rm m} d\omega$ resolution $\omega = \bigcup_i \omega_i$ Up-scaling $\begin{cases} \boldsymbol{\sigma}_{\mathrm{M}} = \frac{1}{V(\omega)} \int_{\omega} \boldsymbol{\sigma}_{\mathrm{m}} d\omega \\ \boldsymbol{q}_{\mathrm{M}} = \frac{1}{V(\omega)} \int_{\omega} \boldsymbol{q}_{\mathrm{m}} d\omega \end{cases}$ $\mathbb{C}_{\mathrm{M}} = \frac{\partial \boldsymbol{\sigma}_{\mathrm{M}}}{\partial \boldsymbol{u}_{\mathrm{M}} \otimes \boldsymbol{\nabla}_{\mathrm{M}}} \qquad \& \quad \boldsymbol{\alpha}_{\mathrm{M}} : \mathbb{C}_{\mathrm{M}} = -\frac{\partial \boldsymbol{\sigma}_{\mathrm{M}}}{\partial \boldsymbol{\vartheta}_{\mathrm{M}}}$ $\kappa_{\rm M} = -\frac{\partial q_{\rm M}}{\partial \nabla u^{9} u}$ $\rho_{\rm M} C_{\nu \rm M} = \frac{1}{V(\omega)} \int \rho_{\rm m} C_{\nu \rm m} dV$ Consistency — Satisfied by periodic boundary conditions Beginning April 2019 CM3 research projects 101

Quality factor

- Micro-resonators
 - Temperature changes with compression/traction
 - Energy dissipation
- Eigen values problem
 - Governing equations



- $\begin{bmatrix} \mathbf{M} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} \end{bmatrix} \begin{bmatrix} \ddot{\mathbf{u}} \\ \ddot{\boldsymbol{\vartheta}} \end{bmatrix} + \begin{bmatrix} \mathbf{0} & \mathbf{0} \\ \mathbf{D}_{u\vartheta}(\boldsymbol{\theta}) & \mathbf{D}_{\vartheta\vartheta} \end{bmatrix} \begin{bmatrix} \dot{\mathbf{u}} \\ \dot{\boldsymbol{\vartheta}} \end{bmatrix} + \begin{bmatrix} \mathbf{K}_{uu}(\boldsymbol{\theta}) & \mathbf{K}_{u\vartheta}(\boldsymbol{\theta}) \\ \mathbf{0} & \mathbf{K}_{\vartheta\vartheta}(\boldsymbol{\theta}) \end{bmatrix} \begin{bmatrix} \mathbf{u} \\ \boldsymbol{\vartheta} \end{bmatrix} = \begin{bmatrix} F_{u} \\ F_{\vartheta} \end{bmatrix}$
- Free vibrating problem

$$\begin{bmatrix} \mathbf{u}(t) \\ \boldsymbol{\vartheta}(t) \end{bmatrix} = \begin{bmatrix} \mathbf{u}_{\mathbf{0}} \\ \boldsymbol{\vartheta}_{\mathbf{0}} \end{bmatrix} e^{i\omega t}$$

$$\begin{array}{c|cccc} & -\mathbf{K}_{\mathrm{uu}}(\boldsymbol{\theta}) & -\mathbf{K}_{\mathrm{u}\vartheta}(\boldsymbol{\theta}) & \mathbf{0} \\ \mathbf{0} & -\mathbf{K}_{\vartheta\vartheta}(\boldsymbol{\theta}) & \mathbf{0} \\ \mathbf{0} & \mathbf{0} & \mathbf{I} \end{array} \begin{bmatrix} \mathbf{u} \\ \boldsymbol{\vartheta} \\ \dot{\mathbf{u}} \end{bmatrix} = i\omega \begin{bmatrix} \mathbf{0} & \mathbf{0} & \mathbf{M} \\ \mathbf{D}_{\vartheta\mathrm{u}}(\boldsymbol{\theta}) & \mathbf{D}_{\vartheta\vartheta} & \mathbf{0} \\ \mathbf{I} & \mathbf{0} & \mathbf{0} \end{bmatrix} \begin{bmatrix} \mathbf{u} \\ \boldsymbol{\vartheta} \\ \dot{\mathbf{u}} \end{bmatrix}$$

- Quality factor
 - From the dissipated energy per cycle

•
$$Q^{-1} = \frac{2|\Im\omega|}{\sqrt{(\Im\omega)^2 + (\Re\omega)^2}}$$



- Application of the 3-Scale method to extract the quality factor distribution
 - 3D models readily available
 - The effect of the anchor can be studied



- Surface topology: asperity distribution
 - Upper surface topology by AFM (Atomic Force Microscope) measurements on 2 µmthick poly-silicon films



Deposition temperature [°C]	580	610	630	650
Std deviation [nm]	35.6	60.3	90.7	88.3

AFM data provided by IMT Bucharest, Rodica Voicu, Angela Baracu, Raluca Muller



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Accounting for roughness Second-order homogenization $\boldsymbol{\varepsilon}_{\mathrm{M}}$, $\boldsymbol{\kappa}_{\mathrm{M}}$ $\widetilde{\boldsymbol{n}}_{\mathrm{M}} = \mathbb{C}_{\mathrm{M}_{1}}: \boldsymbol{\varepsilon}_{\mathrm{M}} + \mathbb{C}_{\mathrm{M}_{2}}: \boldsymbol{\kappa}_{\mathrm{M}}$ $\widetilde{\boldsymbol{m}}_{\mathrm{M}} = \mathbb{C}_{\mathrm{M}_3}: \boldsymbol{\varepsilon}_{\mathrm{M}} + \mathbb{C}_{\mathrm{M}_4}: \boldsymbol{\kappa}_{\mathrm{M}}$ Stochastic homogenization $\boldsymbol{\omega} = \cup_i \boldsymbol{\omega}_i$ Several SVE realizations For each SVE $\omega_i = \bigcup_i \omega_i$ The density per unit area is now non-constant • $\mathbb{C}_{\mathbf{m}^{i}}$ 'mⁱ $\mathbb{C}_{\mathrm{m}^i} \ \forall i$ Computational

 $\omega_i = \bigcup_i \omega_i$

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 $\widetilde{m}_{\mathrm{M}}, \mathbb{C}_{\mathrm{M}_3}, \mathbb{C}_{\mathrm{M}_4}$ ----- Meso-scale BVP resolution $\mathbb{C}_{\mathsf{M}_{1}^{j}},\mathbb{C}_{\mathsf{M}_{2}^{j}},\mathbb{C}_{\mathsf{M}_{3}^{j}},\mathbb{C}_{\mathsf{M}_{4}^{j}}$ U_Mj homogenization Samples of the mesoscale homogenized $\mathbb{C}_{\mathbf{m}^{i}}$ elasticity matrix U_M & $\overline{\rho}_{M^{j}}$ density $\overline{\rho}_{M}$ Beginning 105 April 2019 CM3 research projects

 $\widetilde{\boldsymbol{n}}_{M}, \mathbb{C}_{M_{1}}, \mathbb{C}_{M_{2}}$

• Accounting for roughness

- Cantilever of 8 x 3 x $t \,\mu m^3$ deposited at 610 °C

Flat SVEs (no roughness) - F Rough SVEs (Polysilicon film deposited at 610 °C) - R Grain orientation following XRD measurements – Si_{pref} Grain orientation uniformly distributed – Si_{uni} Reference isotropic material – Iso







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- Application to robust design
 - Determination of probabilistic meso-scale properties
 - Propagate uncertainties to higher scale
 - Vibro-meter sensors:
 - Uncertainties in resonance frequency / Q factor

3SMVIB MNT.ERA-NET project

- Open-Engineering, V2i, ULiège (Belgium)
- Polit. Warszawska (Poland)
- IMT, Univ. Cluj-Napoca (Romania)
- Publications (doi)
 - <u>10.1002/nme.5452</u>
 - <u>10.1016/j.cma.2016.07.042</u>
 - <u>10.1016/j.cma.2015.05.019</u>



Computational & Multiscale Mechanics of Materials





DG-Based (Multi-Scale) Fracture

The research has been funded by the Belgian National Fund for Education at the Research in Industry and Farming. SIMUCOMP The research has been funded by the Walloon Region under the agreement no 1017232 (CT-EUC 2010-10-12) in the context of the ERA-NET +, Matera + framework.

The research has been funded by the Walloon Region under the agreement no.7581-MRIPF in the context of the 16th MECATECH call.



April 2019 CM3 research projects
DG-Based Fracture

• Hybrid DG/cohesive law formulation

- Discontinuous Galerkin method
 - Finite-element discretization
 - Same **discontinuous** polynomial approximations for the
 - **Test** functions φ_h and
 - **Trial** functions $\delta \varphi$



 $(a-1)^{-}(a-1)^{+}(a)^{-}(a)^{+}(a+1)^{-}(a+1)^{+}$

- Can easily be combined with a cohesive law for fracture analyses
 - Interface elements already exist
 - Easy to shift from un-fractured to fractured states
 - Remains accurate before
 fracture onset (DG formulation)
 - Efficient // implementation
- Publications (doi)
 - <u>10.1016/j.cma.2010.08.014</u>





Beginning

DG-Based Multi-Scale Fracture

• Multi-scale modeling

- 2 problems are solved concurrently
 - The macro-scale problem
 - The meso-scale problem (on a meso-scale Volume Element)



For meso-scale volume elements embedding crack propagation

 $L_{\text{macro}} >> L_{\text{VE}}$? L_{micro}

For accuracy: Size of the mesoscale volume element smaller than the characteristic length of the macro-scale loading The crack induces a loss of statistical representativeness

• Should recover consistency lost due to the discontinuity



- Micro-Meso fracture model for intra-laminar failure
 - Epoxy-CF (60%), transverse loading _
 - 3 stages captured





Beginning

- Micro-Meso fracture model for intra-laminar failure (2)
 - Scale transition after softening onset
 - Should not depend on the RVE size
 - Extraction of the meso-scale TSL $(\bar{t}_M \text{ vs. } \Delta_M)$







- SIMUCOMP ERA-NET project
 - e-Xstream, CENAERO, ULiège (Belgium)
 - IMDEA Materials (Spain)
 - CRP Henri-Tudor (Luxemburg)
- Publication (doi)
 - <u>10.1016/j.engfracmech.2013.03.018</u>



< Beginning

DG-Based Dynamic Fracture



- Capture triaxiality effects: Cohesive Band Model (CBM)
 - Introduction of a uniform band of given thickness $h_{
 m b}$ [Remmers et al. 2013]





- 1. Bulk stress σ using non-local damage law
- 2. Compute a "band" deformation gradient

$$\mathbf{F}_{\mathrm{b}} = \mathbf{F} + \frac{\llbracket \boldsymbol{u} \rrbracket \otimes \boldsymbol{N}}{h_{\mathrm{b}}} + \frac{1}{2} \boldsymbol{\nabla}_{T} \llbracket \boldsymbol{u} \rrbracket$$

- 3. Band stress σ_b using the (local) damage law
- 4. Recover traction forces $t(\llbracket u \rrbracket, F) = \sigma_b$. *n*
- The cohesive band thickness
 - Evaluated to ensure energy consistency
 - Same dissipated energy as with a damage model



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Band

Bulk

 $\mathbf{F}_{\rm b}, \boldsymbol{\sigma}_{\rm b}$

F, **σ**



Beginning

DG-Based elastic damage to crack transition

• Slit plate



DG-Based elastic damage to crack transition



DG-Based elastic damage to crack transition

Comparison with phase field

- Single edge notched specimen [Miehe et al. 2010]
 - Calibration of damage and CBM parameters with 1D case [Leclerc et al. 2018]



Compact Tension Specimen:

Non-Local damage law combined to cohesive band model improves accuracy





- MRIPF MECATECH project
 - GDTech, UCL, FZ, MECAR, Capital People (Belgium)
- Publication (doi)
 - <u>10.1002/nme.5618</u>
 - <u>10.1016/j.cma.2014.06.031</u>



Computational & Multiscale Mechanics of Materials





Non-local Gurson damage model to crack transition

The research has been funded by the Walloon Region under the agreement no.7581-MRIPF in the context of the 16th MECATECH call.



April 2019 CM3 research projects

• Objective:

- To develop high fidelity numerical methods for ductile failure
- Numerical approach:
 - Combination of 2 complementary methods in a single finite element framework:
 - continuous (damage model)
 - + transition to
 - discontinuous (cohesive band model including triaxiality / strain rate effects)



- Material changes represented via internal variables
 - Constitutive law $\sigma(\varepsilon; Z(t'))$
 - Internal variables Z(t')
 - Different models
 - Lemaitre-Chaboche (degraded properties)
 - Gurson model (yield surface in terms of porosity f)
- Model implementation:
 - Local form
 - Mesh dependency
 - Requires non-local form [Bažant 1988]
 - Introduction of characteristic length l_c
 - Weighted average: $\tilde{Z}(\mathbf{x}) = \int_{V_c} W(\mathbf{y}; \mathbf{x}, l_c) Z(\mathbf{y}) d\mathbf{y}$
 - Implicit form [Peerlings et al. 1998]
 - New degrees of freedom: *Ž*
 - New Helmholtz-type equations: $\tilde{Z} l_c^2 \Delta \tilde{Z} = Z$







The numerical results change without convergence

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Beginning

- Hyperelastic-based formulation
 - Multiplicative decomposition $\mathbf{F} = \mathbf{F}^{e} \cdot \mathbf{F}^{p}, \ \mathbf{C}^{e} = \mathbf{F}^{e^{T}} \cdot \mathbf{F}^{e}, \ J^{e} = \det(\mathbf{F}^{e})$
 - Stress tensor definition
 - Elastic potential $\psi(\mathbf{C}^{e})$
 - First Piola-Kirchhoff stress tensor

$$\mathbf{P} = 2\mathbf{F}^{\mathrm{e}} \cdot \frac{\partial \psi(\mathbf{C}^{\mathrm{e}})}{\partial \mathbf{C}^{\mathrm{e}}} \cdot \mathbf{F}^{\mathrm{p}^{-T}}$$

- Kirchhoff stress tensors
 - In current configuration

$$\boldsymbol{\kappa} = \mathbf{P} \cdot \mathbf{F}^{T} = 2\mathbf{F}^{e} \cdot \frac{\partial \psi(\mathbf{C}^{e})}{\partial \mathbf{C}^{e}} \cdot \mathbf{F}^{e^{T}}$$

- In co-rotational space

$$\boldsymbol{\tau} = \mathbf{C}^{\mathrm{e}} \cdot \mathbf{F}^{\mathrm{e}^{-1}} \boldsymbol{\kappa} \cdot \mathbf{F}^{\mathrm{e}^{-T}} = 2\mathbf{C}^{\mathrm{e}} \cdot \frac{\partial \psi(\mathbf{C}^{\mathrm{e}})}{\partial \mathbf{C}^{\mathrm{e}}}$$

- Logarithmic deformation
 - Elastic potential ψ :

p

$$\psi(\mathbf{C}^{\mathrm{e}}) = \frac{K}{2} \ln^2(J^e) + \frac{G}{4} (\ln(\mathbf{C}^{\mathrm{e}}))^{\mathrm{dev}} : (\ln(\mathbf{C}^{\mathrm{e}}))^{\mathrm{dev}}$$

Stress tensor in co-rotational space

$$\boldsymbol{\tau} = \underbrace{K \ln(J^e)}_{I} \mathbf{I} + G(\ln(\mathbf{C}^e))^{dev}$$





- Porous plasticity (or Gurson) approach
 - Competition between 2 plastic modes:



- Hybrid DG model: use of a Cohesive Band Model (CBM)
 - Principles
 - Substitute TSL of CZM by the behavior of a uniform band of thickness h_b [Remmers et al. 2013]



- Localization criterion
 - Thomason: $\mathbf{N} \cdot \boldsymbol{\tau} \cdot \mathbf{N} C_l^f \tau_y \ge 0$
- Methodology [Leclerc et al. 2018]
 - 1. Compute a band strain tensor $\mathbf{F}_{b} = \mathbf{F} + \frac{\llbracket \mathbf{u} \rrbracket \otimes \mathbf{N}}{h_{b}} + \frac{1}{2} \nabla_{T} \llbracket \mathbf{u} \rrbracket$
 - 2. Compute a band stress tensor $\sigma_b(F_b; Z(\tau))$ using the same CDM as bulk elements
 - 3. Recover a surface traction $t(\llbracket u \rrbracket, F) = \sigma_b. n$
- What is the effect of $h_{\rm b}$ (band thickness)
 - Recover the fracture energy



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700

600

[Besson 2003]

[Huespe 2012]

 $l_{
m mesh} = 170~\mu{
m m}$

 $l_{\rm mesh} = 112 \,\mu{
m m}$ $l_{\rm mesh} = 75 \ \mu {\rm m}$ 9

Comparison with literature [Huespe2012,Besson2003]





• Notched round bar



- MRIPF MECATECH project
 - GDTech, UCL, FZ, MECAR, Capital People (Belgium)
- Publication (doi)
 - <u>10.1002/nme.5618</u>



Computational & Multiscale Mechanics of Materials





Stochastic Multi-Scale Fracture of Polycrystalline Films

Robust design of MEMS: Financial support from F. R. S. - F. N. R. S. under the project number FRFC 2.4508.11



April 2019 CM3 research projects

• Multi-scale modeling



- The macro-scale problem
- The meso-scale problem (on a meso-scale Volume Element)



• For meso-scale volume elements not several orders larger than the microstructure size and embedding crack propagations

For accuracy: Size of the mesoscale volume element smaller than the characteristic length of the macro-scale loading Meso-scale volume element no longer statistically representative:

- Stochastic Volume Elements
- Should recover consistency lost due to the discontinuity



 $L_{\text{macro}} >> L_{\text{VF}} \sim ? L_{\text{micro}}$

Micro-scale model: Silicon crystal ۲

Different fracture strengths and critical energy release rates _









Define a "continuous" strength mapping







April 2019 CM3 research projects

E L3 1.2

 σ_{C}

0.9

0.8

t

 σ_{c}

θ

 G_C

• Micro-scale model: Polycrsytalline films

 $\Delta(+)$

inter-granular

fracture

- <u>Discontinuous Galerkin method</u>
- Extrinsic cohesive law
- Intra/Inter granular fracture
- Accounts for interface orientation

intra-granular fracture

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+)

 (\pm)

Beginning

2

φ

1

0 0

 $\Delta_c \Delta$

- Stochastic micro-scale to meso-scale model
 - <u>Several SVE realizations (random grain orientation)</u>
 - <u>Extraction of consistent meso-scale cohesive laws</u>
 - \bar{t}_M vs. Δ_M

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- for each SVE sample
- Resulting meso-scale cohesive law distribution



- Macro-scale simulation
 - Finite element model nonconforming to the grains
 - Use homogenized (random) mesoscale cohesive laws as input



- Collaboration for experiments – UcL (T. Pardoen, J.-P Raskin)
- Publications
 - <u>10.1007/s00466-014-1083-4</u>





April 2019 CM3 research projects

Computational & Multiscale Mechanics of Materials





Smart Composite Materials

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Smart Composite Materials

- Electro-thermo-mechanical coupling
 - Finite field variation formulation
 - Strong coupling



Conservation of electric charge

 $\begin{aligned} \mathbf{J}_{e} \cdot \mathbf{\nabla}_{0} &= 0\\ \mathbf{J}_{e} &= \mathbf{J}_{e}(\mathbf{F}, \mathbf{\nabla}_{\mathbf{0}} V, V, \mathbf{\nabla}_{\mathbf{0}} \vartheta, \vartheta; \mathbf{Z}) \end{aligned}$

Conservation of energy

 $\rho C_{v} \dot{\vartheta} - \mathcal{D} + \mathbf{J}_{y} \cdot \nabla_{0} = 0$ $\mathbf{J}_{y} = \mathbf{q} + V \mathbf{J}_{e}$ $\mathbf{q} = \mathbf{q}(\mathbf{F}, V, \nabla_{0} \vartheta, \vartheta; \mathbf{Z})$

Conservation of momentum balance

$$\mathbf{P} \cdot \nabla_0 = 0$$

$$\mathbf{P} = \mathbf{P}(\mathbf{F}, \vartheta; \mathbf{Z})$$

$$\mathcal{D} = \beta \dot{p}\tau + \vartheta \frac{\partial \dot{W}^{\text{el}}}{\partial \vartheta}$$

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Beginning

- Two-way electro-thermal coupling
 - Seebeck coefficient α
 - Finite strain conductivities $\mathbf{K}(V, \vartheta) = \mathbf{F}^{-1} \cdot \mathbf{k}(V, \vartheta) \cdot \mathbf{F}^{-T} \mathbf{J} \ \& \ \mathbf{L}(V, \vartheta) = \mathbf{F}^{-1} \cdot \mathbf{l}(V, \vartheta) \cdot \mathbf{F}^{-T} \mathbf{J}$



- The coefficients matrix $\mathbf{Z}(\mathbf{F}, f_V, f_{\vartheta})$ is symmetric and definite positive



Beginning

Smart Composite Materials

- Thermo-mechanical shape memory polymer
 - Deformations above glass transition temperature ϑ_g (1)
 - Fixed once cooled down below ϑ_g (2 & 3)
 - Recovery once heated up (4)

Elasto-visco-plastic model constitutive behavior

- Different mechanisms (α)
 - Multiplicative decomposition $\mathbf{F}^{(\alpha)} = \mathbf{F}^{e(\alpha)} \mathbf{F}^{p(\alpha)}$
 - Free energy

$$\psi = \sum_{\alpha} \psi^{(\alpha)} \left(\mathbf{C}^{\mathrm{e}^{(\alpha)}}, \vartheta \right)$$

• Thermo-visco-plasticity

$$\tau^{(\alpha)} = \mathcal{T}\left(\mathbf{C}^{\mathrm{e}(\alpha)}, \mathbf{F}^{\mathrm{p}(\alpha)}, \dot{p}^{(\alpha)}, \vartheta, \xi^{(\alpha)}\right)$$

Stress and dissipation

$$\begin{cases} \mathbf{P} = \mathbf{P} \Big(\mathbf{F}, \vartheta; \mathbf{F}^{\mathbf{p}(\alpha)}, p^{(\alpha)}, \xi^{(\alpha)} \Big) \\ \mathcal{D} = \beta \dot{p}^{(\alpha)} \tau^{(\alpha)} \end{cases}$$





[V. Srivastav et. al, 2010]



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 σ^{\prime}

Elasto-visco-plastic behavior of thermo-mechanical shape memory polymer



Smart Composite Materials

- Recovery of a shape memory composite unit cell
 - Carbon Fiber reinforced SMP
 - Shape memory effect triggered by Joule effect
 - Test with compressive force recovery:
 - #1: Compression deformation obtained above ϑ_g
 - #2: Fixation of the deformation above ϑ_g
 - #3: Reheat above ϑ_g at constant deformation:
 - → recovery force, the cell wants to expend
 - #4: Release deformation/stress







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Beginning

- Recovery of a shape memory composite unit cell
 - Carbon Fiber reinforced SMP
 - Triggered by Joule effecy





Discontinuous Galerkin implementation

- Finite-element discretization
- Same discontinuous polynomial approximations for the
 - **Test** functions φ_h and
 - **Trial** functions $\delta \varphi$



- Publication (doi)
 - <u>10.1007/s11012-017-0743-9</u>
 - <u>10.1016/j.jcp.2017.07.028</u>





Computational & Multiscale Mechanics of Materials









April 2019 CM3 research projects

Multi-Scale Modeling of Nano-Crystal Grain Boundary Sliding



Crystal plasticity characterization by nano-indentation



Multi-Scale Modeling of Nano-Crystal Grain Boundary Sliding

- Grain size effect
 - Competition between inter-intra granular



Grain size: 3.28 nm

- Effect of nano-voids in the grain boundaries
 - Different deformation mechanism
 - Lower yield stress
- Collaboration
 - EC Nantes, Univ. of Vermont, Oxford
- Publications
 - <u>10.1016/j.commatsci.2014.03.070</u>
 - <u>10.1016/j.actamat.2013.10.056</u>
 - <u>10.1016/j.jmps.2013.04.009</u>





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Beginning

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Stochastic Multi-Scale Model to Predict MEMS Stiction

3SMVIB: The research has been funded by the Walloon Region under the agreement no 1117477 (CT-INT 2011-11-14) in the context of the ERA-NET MNT framework.

The research has been funded by the Belgian National Fund for Education at the Research in Industry and Farming.



April 2019 CM3 research projects
• Stiction (adhesion of MEMS)

- Different physics at the different scales
- Elastic or Elasto-plastic behaviors
- Due to van der Waals (dry environment) and/or capillary (humid environment) forces
- Requires surfaces topology knowledge (AFM measures)
 - Subject to uncertainties





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- Deterministic multi-scale models for van der Waals forces
 - Extraction of meso-scale adhesive-forces
 - Using statistical representations of the rough surface (average solution)
 - Account for induced elasto-plasticity (cyclic loading)



- New multi-scale models with capillary effect
 - Extraction of meso-scale adhesive-forces from a single surface measurement
 - Depends on the surface sample measurement location
 - Motivates the development of a stochastic multi-scale method



• Stochastic multi-scale model: From the AFM to virtual surfaces

Enforce statistical moments with maximum entropy method



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• Stochastic multi-scale model: Evaluate meso-scale surface forces



• Stochastic multi-scale model: Stochastic model of meso-scale adhesion forces



• Stochastic multi-scale model: Stochastic MEMS stiction analyzes



Application to robust design

- Determination of probabilistic meso-scale properties
- Propagate uncertainties to higher scale
- Vibro-meter sensors:
 - Uncertainties in stiction risk

• 3SMVIB MNT.ERA-NET project

- Open-Engineering, V2i, ULiège (Belgium)
- Polit. Warszawska (Poland)
- IMT, Univ. Cluj-Napoca (Romania)

• FNRS-FRIA fellowship

- Publications (doi)
 - <u>10.1109/JMEMS.2018.2797133</u>
 - <u>10.1016/j.triboint.2016.10.007</u>
 - <u>10.1007/978-3-319-42228-2_1</u>
 - <u>10.1016/j.cam.2015.02.022</u>
 - <u>10.1016/j.triboint.2012.08.003</u>
 - 10.1007/978-1-4614-4436-7_11
 - <u>10.1109/JMEMS.2011.2153823</u>
 - <u>10.1063/1.3260248</u>



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